

# The Spectrum Distribution of Transport Operator with Abstract Boundary Conditions in Slab Geometry

HONGXING WU

Shangrao Normal University  
Department of Mathematics  
Shangrao Jiangxi 334001  
CHINA  
jxsruwhx@163.com

SHENGHUA WANG

Shangrao Normal University  
Department of Mathematics  
Shangrao Jiangxi 334001  
CHINA  
jxsruwsh@163.com

DENGBIN YUAN

Shangrao Normal University  
Department of Mathematics  
Shangrao Jiangxi 334001  
CHINA  
yuandengbin@163.com

**Abstract:** In this paper, transport equation with continuous energy, nonhomogeneous medium and abstract boundary conditions is studied in slab geometry. It is to prove that  $\|K(\lambda I - B_H)^{-1}K\|$  ( $|\operatorname{Im}\lambda| \rightarrow +\infty$ ) is bounded in the trip  $\Gamma_\varepsilon$ , and the spectrum of transport operator  $A_H$  consists of only finite isolated eigenvalues with a finite algebraic multiplicities in trip  $\Gamma_\varepsilon$ . The main methods rely on operators theory, resolvent operators and comparison operators approach.

**Key-Words:** Transport operator; abstract boundary condition; boundedness; isolated eigenvalues.

## 1 Related Knowledge

In this paper, we are concerned with the transport equation with continuous energy, nonhomogeneous medium and abstract boundary conditions in slab geometry. The specific model is as follow

$$\begin{aligned} & \frac{\partial \psi(x, v, \mu, t)}{\partial t} = \\ & - \mu \frac{\partial \psi(x, v, \mu, t)}{\partial x} - \sigma(x, v) \psi(x, v, \mu, t) \quad (1) \\ & + \int_E dv' \int_{-1}^1 k(x, v, \mu, v', \mu') \psi(x, v', \mu', t) d\mu', \end{aligned}$$

with the initial condition

$$\psi(x, \mu, v, 0) = \psi_0(x, \mu, v), \quad (2)$$

where the function  $\psi(x, v, \mu)$  represents the number density of gas particles having the position  $x$ , the particle velocity  $v$  and the direction cosine of propagation  $\mu$ . here  $x \in [-a, a]$  for a parameter  $0 < a < +\infty$ ,  $v, v' \in E = [v_m, v_M]$ ,  $0 < v_m < v_M < +\infty$ , and the  $v_m$  and  $v_M$  are called, respectively, minimum velocity and maximum velocity, and  $\mu, \mu' \in [-1, 1]$ , the function  $\sigma(\cdot, \cdot)$  is called the collision frequency, and the function  $k(\cdot, \cdot, \cdot, \cdot, \cdot)$  is called the scattering kernel. The abstract boundary conditions are modeled by

$$\psi^i = H\psi^0, \quad (3)$$

here,  $H$  is a linear operator in boundary space.

Since Lehner and Wing made some creating work in [1] in 1950's, the research of spectral distribution

of the transport equation have been interesting topic in mathematics, physics, biology and sociology. Latrach and Dehici [2] investigated some spectral properties of time-dependent anisotropic transport equation with periodic and perfecting boundary conditions, using the perturbation theory of strongly continuous semigroups. In fact, let  $X$  be a Banach space, and the streaming operator  $B$  generates a  $C_0$  semigroup  $(U(t)_{t \geq 0})$ . It is well known that if  $K \in \mathcal{L}(X)$  is bounded linear operators, then  $B + K$  generates a strongly continuous semigroup  $(V(t)_{t \geq 0})$ , where

$$V(t) = \sum_{j=0}^{n-1} U_j(t) + R_n(t), \quad (4)$$

where  $U_0(t) = U(t)$ , and

$$U_i(t) = \int_0^t U(s) K U_{j-1}(t-s) ds, \quad j = 1, 2, \dots, \quad (5)$$

and the remainder term  $R_n(t)$  can be expressed by

$$\begin{aligned} R_n(t) &= \sum_{j=n}^{+\infty} U_j(t) \\ &= \int_{t_1 + \dots + t_n \leq t, t_i \geq 0} U(t_1) K U(t_2) K \dots U(t_n) \\ &\quad \times K V(t - t_1 - \dots - t_n) dt_1 \dots dt_n, \quad (6) \end{aligned}$$

where if  $n = 2$ , we can get

$$\begin{aligned} R_2(t) &= \int_{t_1 + t_2 \leq t, t_1 \geq 0, t_2 \geq 0} U(t_1) K U(t_2) \\ &\quad \times K V(t - t_1 - t_2) dt_1 dt_2. \quad (7) \end{aligned}$$

The above method was named by semigroup perturbation approach, and this approach was used by many authors (see, e.g., [3]-[7]). Some authors developed the perturbation technique to the essential spectral radius of transport operators (see, e.g., [8]-[12]).

Recently, Wang and Ma in [13] discussed the transport operator of anisotropic continuous energy and homogeneous with periodic boundary conditions in slab geometry in  $L_2$  space. They proved that the streaming operator  $B$  generates a  $C_0$  semigroup  $(U(t)_{t \geq 0})$ , the transport operator  $A$  generates a  $C_0$  semigroup, and the second-order remained term  $R_2(t)$  of the Dyson-Phillips expansion (4) of the  $C_0$  semigroup is compact in  $L_2$  space. Hence the spectra of the transport operator in some vertical strip  $\Gamma$  consists only of finite many isolated eigenvalues that has a finite algebraic multiplicity. Wang and Wu in [14] discussed the transport operator with anisotropic continuous energy and nonhomogeneous with general boundary conditions in slab geometry in  $L_p(1 \leq p < \infty)$  space. They proved that the streaming operator  $B$  generates a  $C_0$  semigroup  $(U(t)_{t \geq 0})$ , where  $U(t)$  is of the form

$$\begin{aligned}
 U(t)\varphi(x, v, u) &= \sum_{n \geq 0} \alpha^{2n} \\
 &\times \exp\left(-\frac{1}{|\mu|}\left(2n \int_{-a}^a + \text{sgn}(\mu) \int_{x'}^x\right)\sigma(\xi, v) d\xi\right) \\
 &\times \varphi(\text{sgn}(\mu)4na + x - \mu t, v, \mu) \\
 &\times \chi_{[(\text{sgn}(\mu)x+(4n-1)a)/|\mu|, (\text{sgn}(\mu)x+(4n+1)a)/|\mu|]}(t) \\
 &+ \sum_{n \geq 0} \alpha^{2n+1} \exp\left(-\frac{2n}{|\mu|} \int_{-a}^a \sigma(\xi, v) d\xi\right) \quad (8) \\
 &\times \exp\left(-\frac{1}{|\mu|} \text{sgn}(\mu)\left(\int_{-a}^x + \int_{-a}^{x'}\right)\sigma(\xi, v) d\xi\right) \\
 &\times \varphi(-\text{sgn}(\mu)(4n+2)a - x + \mu t, v, -\mu) \\
 &\times \chi_{[(\text{sgn}(\mu)x+(4n+1)a)/|\mu|, (\text{sgn}(\mu)x+(4n+3)a)/|\mu|]}(t),
 \end{aligned}$$

the transport operator  $A$  generates a  $C_0$  semigroup, and the second-order remained term  $R_2(t)$  of the Dyson-Phillips expansion of the semigroup is compact in  $L_p(1 < p < \infty)$  space and weakly compact in  $L_1$  space. It is similar to the result of [13].

It is well-known that if the transport equation with the specific boundary conditions, or abstract boundary conditions, then the bounded perturbation methods will fail. This is because the boundary operator is a unbounded linear operator. So we have to use the resolvent analysis approach to study the transport equation. Latrach and Megdiche in [15] discussed the transport equation with anisotropic and abstract boundary conditions in slab geometry. Under some assumption that, for  $r \in [0, 1)$

$$\lim_{|\Im \lambda| \rightarrow +\infty} |\Im \lambda|^r \|K(\lambda I - B)^{-1}K\| = 0, \quad (9)$$

uniformly on some vertical strip, they derived various descriptions of the large time behavior of solutions. Latrach et al. in [16] discussed the transport equation with reentry boundary conditions in slab geometry, they derived conditions that ensure the compactness of the remainder term  $R_n(t)$  for some integer  $n$ , and got the large time asymptotic behavior of the solution to the one-dimensional transport equation. Lately, some authors discussed the transport equation with anisotropic continuous energy and homogeneous in slab geometry, and obtained essential spectrum and isolated spectrum of the transport equation (see, e.g., [17]-[24], [30]-[32]).

In the past years, some authors described the time asymptotic behavior of the solution of a one-velocity transport operator without restriction on the initial data in sphere (see, e.g., [25, 26]). Of course, there are some progresses about the spectral of bizarre transport equation (see, e.g., [27, 28]). The spectral analysis of transport operator in growing cell population (see, e.g., [33-35]). Recent, Abdelmoumen et al. in [29] discussed the transport operator with anisotropic in sphere, and described the large time behavior of solutions to an abstract Cauchy problem under some assumptions. They proved that there exists an integer  $m_0$  and  $r_0 \in [0, 1)$  such that

$$\|\Im \lambda\|^{r_0} \|[(\lambda I - B)^{-1}K]^{m_0}\|, \quad (10)$$

is bounded uniformly in some vertical strip. A question is what spectral distribution in slab geometry is under the above condition. In this paper, we will discuss, in  $L_p(1 \leq p < +\infty)$  space, the transport equation with continuous energy nonhomogeneous medium and abstract boundary conditions in slab geometry. We will prove that operator

$$\|\Im \lambda\| \|K(\lambda I - B_H)^{-1}K\|, (|\Im \lambda| \rightarrow +\infty), \quad (11)$$

is bounded on a vertical strip  $\Gamma_\varepsilon$ , and the spectrum of transport operator in the strip  $\Gamma_\varepsilon$  is composed of finite many isolated eigenvalues of finite algebraic multiplicities.

Let us introduce some notion and notations, and make precise the function setting of the problem. Let space be

$$X = L_p(D, dx dv d\mu), \quad (12)$$

the norm is defined by

$$\|\psi\|_X = \left(\int_{-a}^a \int_E \int_{-1}^1 |\psi(x, v, \mu)|^p dx dv d\mu\right)^{\frac{1}{p}}, \quad (13)$$

where  $D = [-a, a] \times E \times [-1, 1], p \in [1, +\infty)$ .

We define the following sets representing the incoming and the outgoing boundary of the phase space

$$D^0 = D_1^0 \cup D_2^0 = \{-a\} \times E \times [-1, 0] \cup \{a\} \times E \times [0, 1], \quad (14)$$

$$D^i = D_1^i \cup D_2^i = \{-a\} \times E \times [0, 1] \cup \{a\} \times E \times [-1, 0]. \quad (15)$$

Moreover, we introduce the following boundary spaces

$$\begin{aligned} X^0 &= L_p(D^0, |\mu|dvd\mu) \sim L_p(D_1^0, |\mu|dvd\mu) \\ &\quad \oplus L_p(D_2^0, |\mu|dvd\mu) \\ &= X_1^0 \oplus X_2^0, \end{aligned} \quad (16)$$

$$\begin{aligned} X^i &= L_p(D^i, |\mu|dvd\mu) \sim L_p(D_1^i, |\mu|dvd\mu) \\ &\quad \oplus L_p(D_2^i, |\mu|dvd\mu) \\ &= X_1^i \oplus X_2^i, \end{aligned} \quad (17)$$

endowed with the norm

$$\begin{aligned} \|\varphi^0\|_{X^0} &= \left( \|\varphi_1^0\|_{X_1^0}^p + \|\varphi_2^0\|_{X_2^0}^p \right)^{\frac{1}{p}} \\ &= \left( \int_E dv \int_{-1}^0 |\varphi(-a, v, \mu)|^p |\mu|d\mu \right. \\ &\quad \left. + \int_E dv \int_0^1 |\varphi(a, v, \mu)|^p |\mu|d\mu \right)^{\frac{1}{p}}, \end{aligned} \quad (18)$$

$$\begin{aligned} \|\varphi^i\|_{X^i} &= \left( \|\varphi_1^i\|_{X_1^i}^p + \|\varphi_2^i\|_{X_2^i}^p \right)^{\frac{1}{p}} \\ &= \left( \int_E dv \int_0^1 |\varphi(-a, v, \mu)|^p |\mu|d\mu \right. \\ &\quad \left. + \int_E dv \int_{-1}^0 |\varphi(a, v, \mu)|^p |\mu|d\mu \right)^{\frac{1}{p}}, \end{aligned} \quad (19)$$

where  $\sim$  means the natural identification of the above spaces. We define the streaming operator  $B_H$  by

$$B_H\psi(x, v, \mu) = -\mu \frac{\partial \psi(x, v, \mu)}{\partial x} - \sigma(x, v)\psi(x, v, \mu), \quad (20)$$

where

$$D(B_H) = \left\{ \psi \in X \left| \mu \frac{\partial \psi}{\partial x} \in X, \psi^i = H\psi^0 \right. \right\}, \quad (21)$$

where  $\sigma(x, v)$  is a non-negative and measurable function,  $\psi^0 = (\psi_1^0, \psi_2^0)^\top$ , and  $\psi^i = (\psi_1^i, \psi_2^i)^\top$  with  $\psi_1^0, \psi_2^0, \psi_1^i$  and  $\psi_2^i$  are given by

$$\psi_1^i(v, \mu) = \psi(-a, v, \mu), \quad (22)$$

$$\psi_2^i(v, \mu) = \psi(a, v, \mu), \quad (23)$$

$$\psi_1^0(v, \mu) = \psi(-a, v, \mu), \quad (24)$$

$$\psi_2^0(v, \mu) = \psi(a, v, \mu). \quad (25)$$

Moreover, we define the disturbance operators  $K$  by

$$\begin{aligned} K\psi(x, v, \mu) &= \int_E dv' \int_{-1}^1 \\ &\quad \times k(x, v, \mu, v', \mu')\psi(x, v', \mu')d\mu'. \end{aligned} \quad (26)$$

So, we can define the transport operator  $A_H$  by

$$A_H = B_H + K, \quad D(A_H) = D(B_H). \quad (27)$$

Setting

$$\sigma_0 = \text{essinf}\{\sigma(x, v)\}.$$

Let  $\varphi \in X$  and consider the resolvent equation for  $B_H$

$$(\lambda I - B_H)\psi = \varphi. \quad (28)$$

Thus, for  $\Re\lambda > -\sigma_0$ , the solution of (28) is formally given by

$$\begin{aligned} &\psi(x, v, \mu) \\ &= \psi(-a, v, \mu) \exp\left(\frac{-1}{\mu} \int_{-a}^x (\lambda + \sigma(\xi, v))d\xi\right) \\ &\quad + \frac{1}{\mu} \int_{-a}^x \exp\left(\frac{-1}{\mu} \int_{x'}^x (\lambda + \sigma(\xi, v))d\xi\right) \\ &\quad \times \varphi(x', v, \mu)dx', \quad \mu \in (0, 1), \end{aligned} \quad (29)$$

$$\begin{aligned} &\psi(x, v, \mu) \\ &= \psi(a, v, \mu) \exp\left(\frac{1}{\mu} \int_x^a (\lambda + \sigma(\xi, v))d\xi\right) \\ &\quad - \frac{1}{\mu} \int_x^a \exp\left(\frac{1}{\mu} \int_x^{x'} (\lambda + \sigma(\xi, v))d\xi\right) \\ &\quad \times \varphi(x', v, \mu)dx', \quad \mu \in (-1, 0). \end{aligned} \quad (30)$$

For  $x = \pm a$ , we can get

$$\begin{aligned} &\psi(a, v, \mu) \\ &= \psi(-a, v, \mu) \exp\left(\frac{-1}{\mu} \int_{-a}^a (\lambda + \sigma(\xi, v))d\xi\right) \\ &\quad + \frac{1}{\mu} \int_{-a}^a \exp\left(\frac{-1}{\mu} \int_{x'}^a (\lambda + \sigma(\xi, v))d\xi\right) \\ &\quad \times \varphi(x', v, \mu)dx', \end{aligned} \quad (31)$$

$$\begin{aligned} & \psi(-a, v, \mu) \\ = & \psi(a, v, \mu) \exp\left(\frac{1}{\mu} \int_{-a}^a (\lambda + \sigma(\xi, v)) d\xi\right) \\ & - \frac{1}{\mu} \int_{-a}^a \exp\left(\frac{1}{\mu} \int_{-a}^{x'} (\lambda + \sigma(\xi, v)) d\xi\right) \\ & \times \varphi(x', v, \mu) dx'. \end{aligned} \quad (32)$$

Now, we define operators  $P_\lambda, Q_\lambda, D_\lambda$  and  $E_\lambda$  as follows

$$P_\lambda : X^i \rightarrow X^0; \quad P_\lambda \varphi = (P_\lambda^+ \varphi, P_\lambda^- \varphi), \quad (33)$$

where

$$\begin{aligned} P_\lambda^+ \varphi(a, v, \mu) &= \varphi(-a, v, \mu) \\ &\times \exp\left(\frac{-1}{\mu} \int_{-a}^a (\lambda + \sigma(\xi, v)) d\xi\right), \end{aligned} \quad (34)$$

and

$$\begin{aligned} P_\lambda^- \varphi(-a, v, \mu) &= \varphi(a, v, \mu) \\ &\times \exp\left(\frac{1}{\mu} \int_{-a}^a (\lambda + \sigma(\xi, v)) d\xi\right); \end{aligned} \quad (35)$$

$$Q_\lambda : X^i \rightarrow X; \quad Q_\lambda \varphi = (Q_\lambda^+ \varphi, Q_\lambda^- \varphi), \quad (36)$$

where

$$\begin{aligned} Q_\lambda^+ \varphi(-a, v, \mu) &= \varphi(-a, v, \mu) \\ &\times \exp\left(\frac{-1}{\mu} \int_{-a}^x (\lambda + \sigma(\xi, v)) d\xi\right), \end{aligned} \quad (37)$$

and

$$\begin{aligned} Q_\lambda^- \varphi(a, v, \mu) &= \varphi(a, v, \mu) \\ &\times \exp\left(\frac{1}{\mu} \int_x^a (\lambda + \sigma(\xi, v)) d\xi\right); \end{aligned} \quad (38)$$

$$D_\lambda : X \rightarrow X^0; \quad D_\lambda \varphi = (D_\lambda^+ \varphi, D_\lambda^- \varphi), \quad (39)$$

where

$$\begin{aligned} D_\lambda^+ \varphi(x, v, \mu) &= \frac{1}{\mu} \int_{-a}^a \varphi(x', v, \mu) \\ &\times \exp\left(\frac{-1}{\mu} \int_{x'}^a (\lambda + \sigma(\xi, v)) d\xi\right) dx', \end{aligned} \quad (40)$$

and

$$\begin{aligned} D_\lambda^- \varphi(-a, v, \mu) &= \frac{1}{\mu} \int_{-a}^a \varphi(x', v, \mu) \\ &\times \exp\left(\frac{1}{\mu} \int_{-a}^{x'} (\lambda + \sigma(\xi, v)) d\xi\right) dx'; \end{aligned} \quad (41)$$

$$E_\lambda : X \rightarrow X; \quad E_\lambda \varphi = (E_\lambda^+ \varphi, D_\lambda^- \varphi), \quad (42)$$

where

$$\begin{aligned} E_\lambda^+ \varphi(x, v, \mu) &= \frac{1}{\mu} \int_{-a}^x \varphi(x', v, \mu) \\ &\times \exp\left(\frac{-1}{\mu} \int_{x'}^x (\lambda + \sigma(\xi, v)) d\xi\right) dx, \end{aligned} \quad (43)$$

and

$$\begin{aligned} E_\lambda^- \varphi(-a, v, \mu) &= \frac{1}{\mu} \int_x^a \varphi(x', v, \mu) \\ &\times \exp\left(\frac{1}{\mu} \int_x^{x'} (\lambda + \sigma(\xi, v)) d\xi\right) dx. \end{aligned} \quad (44)$$

We assume that the boundary operator  $H$  satisfies the following condition.

**Assumption  $O_1$ :**  $H : X^0 \rightarrow X^i$ ,

$$H \begin{pmatrix} u_1 \\ u_2 \end{pmatrix} = \begin{pmatrix} 0 & H_{12} \\ H_{21} & 0 \end{pmatrix} \begin{pmatrix} u_1 \\ u_2 \end{pmatrix}. \quad (45)$$

where

$$\begin{cases} H_{12} = \alpha J_1 + \beta L_1 : X_2^0 \rightarrow X_2^i; \\ H_{12} \in L(X_2^0, X_2^i), \end{cases} \quad (46)$$

$$\begin{cases} H_{21} = \alpha J_2 + \beta L_2 : X_1^0 \rightarrow X_2^i; \\ H_{21} \in L(X_1^0, X_2^i), \end{cases} \quad (47)$$

$\alpha, \beta \in R^+, J_1$  and  $J_2$  are compact operators. Moreover

$$L_1 u(-a, v, \mu) = u(a, v, \mu), \quad (48)$$

$$L_2 u(a, v, \mu) = u(-a, v, \mu). \quad (49)$$

So, for  $\Re \lambda > -\sigma_0$ , we get

$$\begin{aligned} (\lambda I - B_H)^{-1} &= \chi_{(0,1)}(\mu) R^+(\lambda I, B_H) \\ &+ \chi_{(-1,0)}(\mu) R^-(\lambda I, B_H), \end{aligned} \quad (50)$$

where,

$$\begin{aligned} & R^+(\lambda I, B_H) \\ &= \sum_{n \geq 0} Q_\lambda^+ H_{12} (P_\lambda^+ H_{12})^n D_\lambda^+ + E_\lambda^+, \end{aligned} \quad (51)$$

$$\begin{aligned} & R^-(\lambda I, B_H) \\ &= \sum_{n \geq 0} Q_\lambda^- H_{21} (P_\lambda^- H_{21})^n D_\lambda^- + E_\lambda^-. \end{aligned} \quad (52)$$

**Assumption  $O_2$ :** Operator  $K$  is a regular operator in  $X$ . So it can be approximated in the uniform operator topology by operators. Thus

$$\begin{aligned} K \varphi(x, v, \mu) &= \sum_{i \in I} \int_E dv' \int_{-1}^1 \theta_i(x) f_i(v, \mu) \\ &\times g_i(v', \mu') \varphi(x, v', \mu') d\mu', \end{aligned} \quad (53)$$

where  $\theta_i(\cdot) \in L_\infty([-a, a])$ ,  $f_i(\cdot, \cdot) \in L_1(E \times [-1, 1])$ ,  $g_i(\cdot, \cdot) \in L_\infty(E \times [-1, 1])$ ,  $I$  is finite set.  
Setting

$$\lambda_0 = \begin{cases} -\sigma_0, & \|H\| \leq 1, \\ -\sigma_0 + \frac{1}{2a} \log(\|H\|), & \|H\| > 1. \end{cases} \quad (54)$$

**Lemma 1.** [15] *If the assume  $O_1$  is satisfied, then, for  $\Re\lambda > -\sigma_0$ , we have  $(\lambda I - B_H)^{-1}$  is bounded and*

$$\|(\lambda I - B_H)^{-1}\| \leq \frac{1}{\Re\lambda + \sigma_0}. \quad (55)$$

**Lemma 2.** [5] *If for any  $\varepsilon > 0$ , there exists a  $m \in N$ ,  $\eta$ , such that  $[(\lambda I - B_H)^{-1}K]^m$  is compact, and*

$$\lim_{|\Im\lambda| \rightarrow +\infty} \|[(\lambda I - B_H)^{-1}K]^m\| = 0. \quad (56)$$

*Then, there exists at most finitely many isolated eigenvalues of  $A_H$  in the strip  $\{\lambda \in \mathbb{C}; \Re\lambda \geq \eta + \varepsilon\}$  where  $\eta$  is type of  $C_0$  semigroup generated by streaming operator  $B_H$ , which are of finite algebraic multiplicity.*

## 2 Main Result

In this section, we will give the main results of this paper. Setting

$$\Gamma_\varepsilon = \{\lambda \in C; \Re\lambda \geq -\sigma_0 + \varepsilon\} (\varepsilon > 0). \quad (57)$$

**Theorem 3.** *If assumptions  $O_1$  and  $O_2$  are satisfied, then*

$$|\Im\lambda| \|K(\lambda I - B_H)^{-1}K\|, \quad (58)$$

*is uniformly bounded on  $\Gamma_\varepsilon$ .*

**Proof.** We finish the proof by the following serval steps.

**Step 1.** Because of

$$\begin{aligned} & \|K(\lambda - B_H)^{-1}K\| \\ & \leq \|KE_\lambda^+K\| + \|KE_\lambda^-K\| \\ & + \sum_{n \geq 0} \|KQ_\lambda^+H_{12}(P_\lambda^+H_{12})^nD_\lambda^+K\| \\ & + \sum_{n \geq 0} \|KQ_\lambda^-H_{21}(P_\lambda^-H_{21})^nD_\lambda^-K\|. \end{aligned} \quad (59)$$

So, if we prove (57) is bounded uniformly on  $\Gamma_\varepsilon$ , we only prove

$$|\Im\lambda| \|KE_\lambda^+K\|, \quad (60)$$

$$|\Im\lambda| \|KE_\lambda^-K\|, \quad (61)$$

$$|\Im\lambda| \sum_{n \geq 0} \|KQ_\lambda^+H_{12}(P_\lambda^+H_{12})^nD_\lambda^+K\|, \quad (62)$$

$$|\Im\lambda| \sum_{n \geq 0} \|KQ_\lambda^-H_{21}(P_\lambda^-H_{21})^nD_\lambda^-K\|. \quad (63)$$

are all bounded uniformly on  $\Gamma_\varepsilon$ .

**Step 2.** Prove equation (60) is bounded uniformly on  $\Gamma_\varepsilon$ . For all  $\varphi \in X$ , we get

$$\begin{aligned} E_\lambda^+\varphi(x, v, \mu) &= \frac{1}{\mu} \int_{-a}^x \varphi(x', v, \mu) \\ &\times \exp\left(\frac{-1}{\mu} \int_{x'}^x (\lambda + \sigma(\xi, v))d\xi\right) dx' \\ &= \frac{1}{\mu} \int_{-a}^x \varphi(x', v, \mu) \\ &\times \exp\left(\frac{-1}{\mu} \left[(x - x')\lambda + \int_{x'}^x \lambda + \sigma(\xi, v)d\xi\right]\right) dx'. \end{aligned} \quad (64)$$

The change of  $s = \frac{x-x'}{\mu}$  gives

$$\begin{aligned} E_\lambda^+\varphi(x, v, \mu) &= \int_0^{+\infty} \varphi(x - s\mu, v, \mu)\chi_{(0, \frac{x+a}{\mu})}(s) \\ &\times \exp\left(-\lambda s - \int_{x-s\mu}^x \sigma(\xi, v)d\xi\right) ds. \end{aligned} \quad (65)$$

Now consider the sequence of operators  $E_{\lambda, \varepsilon_n}^+$ , where

$$\begin{aligned} & E_{\lambda, \varepsilon_n}\varphi(x, v, \mu) \\ &= \int_{\varepsilon_n}^{+\infty} \varphi(x - s\mu, v, \mu)\chi_{(0, \frac{x+a}{\mu})}(s) \\ &\times \exp\left(-\lambda s - \int_{x-s\mu}^x \sigma(\xi, v)d\xi\right) ds, \end{aligned} \quad (66)$$

where  $(\varepsilon_n)_{n \in N}$  is a sequence of non-negative real numbers which converge to zero as  $n \rightarrow \infty$ . Clearly, the sequence  $(E_{\lambda, \varepsilon_n})_{n \in N}$  converges to  $E_\lambda^+$ , in the operator topology, uniformly on  $\Gamma_\varepsilon$  as  $n \rightarrow \infty$ . So, it suffices to prove that, for  $\varepsilon > 0$ ,

$$|\Im\lambda| \|KE_{\lambda, \varepsilon}^+K\|,$$

is bounded uniformly on  $\Gamma_\varepsilon$ . Because of

$$\begin{aligned} & KE_{\lambda, \varepsilon}^+K\varphi(x, v, \mu) \\ &= \int_E dv' \int_0^1 d\mu' h(v', \mu') f(v, \mu)\chi_{(0, \frac{x+a}{\mu'})}(s) \\ &\times \exp\left(-\lambda s - \int_{x-s\mu'}^x \sigma(\xi, v)d\xi'\right) \\ &\times \int_\varepsilon^{+\infty} \int_E \int_{-1}^1 \theta(x - s\mu')g(v'', \mu'') \\ &\times \theta(x)\varphi(x - s\mu', v'', \mu'') ds dv'' d\mu''. \end{aligned} \quad (67)$$

Setting  $t = x - \mu's$ , we get

$$\begin{aligned} & KE_{\lambda,\varepsilon}^+ K\varphi(x, v, \mu) \\ = & \theta(x) \int_E dv' \int_{-a}^x d\mu' h\left(v', \frac{x-t}{s}\right) f(v, \mu) \\ & \times \exp\left(-\lambda s - \int_t^x \sigma(\xi, v) d\xi'\right) \\ & \times \int_{\varepsilon}^{+\infty} ds \varphi(t', v'', \mu'') \cdot \chi_{(x-t, +\infty)}(s) \\ & \times \int_E \theta(t) g(v'', \mu'') dv'' \int_{-1}^1 d\mu'' \end{aligned} \quad (68)$$

Putting

$$KE_{\lambda,\varepsilon}^+ K = A_1 A_{\varepsilon} A_2, \quad (69)$$

where  $A_1 : L_p(-a, a) \rightarrow X$ ,

$$A_1 \varphi(x) = \theta(x) f(v, \mu) \varphi(x); \quad (70)$$

$A_2 : L_p(-a, a) \rightarrow X$ ,

$$\begin{aligned} & A_2 \varphi(x, v, \mu) \\ = & \int_E dv \int_{-1}^1 \theta(x) g(v, \mu) \varphi(x, v, \mu) d\mu; \end{aligned} \quad (71)$$

$A_{\varepsilon} : L_p(-a, a) \rightarrow L_p(-a, a)$ ,

$$\begin{aligned} A_{\varepsilon} \varphi(x) &= \int_{-a}^x dt \int_E dv \int_{\varepsilon}^{+\infty} \frac{ds}{s} \\ & \times \exp\left[-\lambda s - \int_t^x \sigma(x, \xi) d\xi\right] \\ & \times h\left(v, \frac{x-t}{s}\right) \varphi(t) \chi_{(x-t, +\infty)}(s); \end{aligned} \quad (72)$$

and  $A_{\varepsilon,n} : L_p(-a, a) \rightarrow L_p(-a, a)$ ,

$$\begin{aligned} A_{\varepsilon,n} \varphi(x) &= \int_{-a}^x dt \int_E dv \int_{\varepsilon}^{+\infty} l_{x-t,v,n}(s) \\ & \times \exp\left[-\left(\lambda + \sigma_0 - \frac{\varepsilon}{2}\right)s\right] ds \\ & \times h\left(v, \frac{x-t}{s}\right) \varphi(t) \chi_{(x-t, +\infty)}(s), \end{aligned} \quad (73)$$

where  $l_{x-t,v,n}(\cdot)$  converges to  $\varphi_{x-t,v}(\cdot)$ .a.e. Because operators  $A_1$  and  $A_2$  are bounded and uniformly on  $\Gamma_{\varepsilon}$ , so, it suffices to prove that, for  $\varepsilon > 0$ ,  $\|\mathfrak{S}\lambda\|^p \|A_{\varepsilon}\|^p$  is bounded uniformly on  $\Gamma_{\varepsilon}$ . According to Lemma 1 of [15], it holds that

$$\|\mathfrak{S}\lambda\|^p \|A_{\varepsilon,n}\|^p, (n \in N), \quad (74)$$

is bounded uniformly on  $\Gamma_{\varepsilon}$ . For all  $\varphi \in L_p(-a, a), n \in N$ , we get

$$\begin{aligned} \|A_{\varepsilon,n} \varphi\|^p &= \int_{-a}^a dx \left| \int_{-a}^x dt \int_E dv \int_{\varepsilon}^{+\infty} \right. \\ & \times \exp\left[-\left(\lambda + \sigma_0 - \frac{\varepsilon}{2}\right)s\right] l_{x-t,v,n}(s) \\ & \times h\left(v, \frac{x-t}{s}\right) \varphi(t) \chi_{(x-t, +\infty)}(s) ds \Big|^p \\ & \leq \int_{-2a}^{2a} dx \left| \int_{-a}^x dt \int_E dv \int_{\varepsilon}^{+\infty} \right. \\ & \times \exp\left[-\left(\lambda + \sigma_0 - \frac{\varepsilon}{2}\right)s\right] l_{x,v,n}(s) \\ & \times h\left(v, \frac{x}{s}\right) \varphi(t) \chi_{(x-t, +\infty)}(s) ds \Big|^p. \end{aligned} \quad (75)$$

The use of the Hölder inequality gives

$$\begin{aligned} & \left| \int_{-a}^x dt \int_E dv \int_{\varepsilon}^{+\infty} \exp\left[-\left(\lambda + \sigma_0 - \frac{\varepsilon}{2}\right)s\right] \right. \\ & \times l_{x,v,n}(s) h\left(v, \frac{x}{s}\right) \varphi(t) \chi_{(x-t, +\infty)}(s) ds \Big|^p \\ & \leq \left[ \int_{-a}^a dt \left| \int_E dv \int_{\varepsilon}^{+\infty} \exp\left[-\left(\lambda + \sigma_0 - \frac{\varepsilon}{2}\right)s\right] \right. \right. \\ & \times l_{x,v,n}(s) h\left(v, \frac{x}{s}\right) ds \Big|^q \int_{-a}^a |\varphi(t)|^p dt \\ & \leq (2a)^{\frac{p}{q}} \left| \int_E dv \int_{\varepsilon}^{+\infty} \exp\left[-\left(\lambda_0 + \sigma - \frac{\varepsilon}{2}\right)s\right] \right. \\ & \times l_{x,v,n}(s) h\left(v, \frac{x}{s}\right) ds \Big|^p \|\varphi\|^p. \end{aligned} \quad (76)$$

Use of the Hölder inequality again, we get

$$\begin{aligned} \|A_{\varepsilon,n}\|^p &\leq (2a)^{\frac{p}{q}} \int_{-2a}^{2a} dx \left| \int_E dv \int_{\varepsilon}^{+\infty} \right. \\ & \times \exp\left[-\left(\lambda + \sigma_0 - \frac{\varepsilon}{2}\right)s\right] \\ & \times l_{x,v,n}(s) h\left(v, \frac{x}{s}\right) \varphi(t) \chi_{(x, +\infty)}(s) ds \Big|^p \\ & \leq (2a)^{\frac{p}{q}} \int_{-2a}^{2a} dx \int_E dv \left| \int_{\varepsilon}^{+\infty} \right. \\ & \times \exp\left[-\left(\lambda + \sigma_0 - \frac{\varepsilon}{2}\right)s\right] \\ & \times l_{x,v,n}(s) h\left(v, \frac{x}{s}\right) \varphi(t) \chi_{(x, +\infty)}(s) ds \Big|^p. \end{aligned} \quad (77)$$

So, it suffices to prove that, for  $\varepsilon > 0$ ,

$$\begin{aligned} & \left| \mathfrak{S}\lambda \right|^p \int_{-2a}^{2a} dx \int_E dv \left| \int_{\varepsilon}^{+\infty} ds \right. \\ & \times \exp\left[-\left(\lambda + \sigma_0 - \frac{\varepsilon}{2}\right)s\right] s l_{x,v,n}(s) \\ & \times h\left(v, \frac{x}{s}\right) \varphi(t) \chi_{(x, +\infty)}(s) \Big|^p, \end{aligned} \quad (78)$$

is bounded uniformly on  $\Gamma_\varepsilon$ . In fact, for  $\forall n \in N$ ,  $x \in (-2a, 2a)$  and  $v \in E$  be fixed, we define

$$W_{x,v}(\cdot) : (\varepsilon, +\infty) \rightarrow R, \tag{79}$$

$$s \mapsto l_{x,v,n}(s)h(v, \frac{x}{s}), \tag{80}$$

where  $l_{x,v,n}(s)$  and  $h(v, \frac{x}{s})$  are simple functions. Setting  $(s_i)_{1 \leq i \leq m}$ , for  $\forall i \in \{1, 2, \dots, m-1\}, s \in [s_i, s_{i+1})$ , we get

$$W_{x,v}(\cdot) = W_{x,v}(s_i), \tag{81}$$

so we can get

$$\begin{aligned} & \int_\varepsilon^{+\infty} \exp(-(\lambda + \sigma_0 - \frac{\varepsilon}{2})s)W_{x,v}(s)ds \\ &= \sum_{i=1}^{m-1} G_{x,v}(s_i) \int_{s_i}^{s_{i+1}} \exp(-(\lambda + \sigma_0 - \frac{\varepsilon}{2})s)ds \\ &= \sum_{i=1}^{m-1} \left( \exp(-(\lambda + \sigma_0 - \frac{\varepsilon}{2})s_i) \right. \\ & \quad \left. - \exp(-(\lambda + \sigma_0 - \frac{\varepsilon}{2})s_{i+1}) \right) \\ & \quad \times \frac{1}{\lambda + \sigma - \frac{\varepsilon}{2}} W_{x,v}(s_i). \end{aligned}$$

So

$$\left| \int_\varepsilon^{+\infty} W_{x,v}(s)ds \right| \leq \frac{2(m-1) \sup |h(\cdot, \cdot)|}{\varepsilon |\Im \lambda|}, \tag{82}$$

and

$$\begin{aligned} & |\Im \lambda|^p \int_{-2a}^{2a} dx \int_E dv \left| \int_\varepsilon^{+\infty} l_{x,v,n}(s)h(v, \frac{x}{s}) \right. \\ & \quad \left. \times \exp(-(\lambda + \sigma_0 - \frac{\varepsilon}{2})s)\varphi(t)\chi_{(x,+\infty)}(s) \right|^p \\ & \leq |\Im \lambda|^p \int_{-2a}^{2a} dx \int_E dv (2(m-1)) \\ & \quad \times \sup |h(\cdot, \cdot)|^p \varepsilon^{-p} |\Im \lambda|^{-p} \\ & \leq 4aM(2(m-1) \sup |h(\cdot, \cdot)|^p \varepsilon^{-p}) |\Im \lambda|^{-p}, \end{aligned} \tag{83}$$

where  $M = v_M - v_m$ . Since

$$4aM(2(m-1) \sup |h(\cdot, \cdot)|^p \varepsilon^{-p}) |\Im \lambda|^{-p},$$

is bounded and uniformly on  $\Gamma_\varepsilon$ . This ends the step 2.

Since the (61) and (60) have the same mode, similarly, we can get the equation (61) is bounded uniformly on  $\Gamma_\varepsilon$ .

**Step 3.** (62) is bounded uniformly on  $\Gamma_\varepsilon$ . Since  $(P_\lambda^+ H_{12})^n$  can be expressed by

$$(P_\lambda^+ H_{12})^n = \sum_{j=1}^{2^n} P_j, \tag{84}$$

where each  $P_j$  is the product of n factors involving both  $\alpha P_\lambda^+ T_1$  and  $\beta P_\lambda^+ L_1$  except the term

$$P_{2^n} = (\beta P_\lambda^+ L_1)^n. \tag{85}$$

So, for  $j \in \{1, 2, \dots, 2^n - 1\}$ , the operator  $T_1$  appears at least once in the expression of  $P_j$ . While

$$\begin{aligned} & \| KQ_\lambda^+ H_{12} P_j D_\lambda^+ K \| \\ & \leq \| KQ_\lambda^+ H_{12} \| \cdot \| P_j D_\lambda^+ K \|, \end{aligned} \tag{86}$$

where  $j \in \{1, 2, \dots, 2^n - 1\}$ , so if we prove the equation (62) is bounded and uniformly on  $\Gamma_\varepsilon$ . We only need to prove

$$|\Im \lambda| \| P_j D_\lambda^+ K \|, \tag{87}$$

$$|\Im \lambda| \| KQ_\lambda^+ H_{12} P_{2^n} D_\lambda^+ K \|, \tag{88}$$

are all bounded uniformly on  $\Gamma_\varepsilon$ , where  $j \in \{1, 2, \dots, 2^n - 1\}$ . In fact, according to the hypotheses, there exists  $k \in \{1, 2, \dots, n-1\}$ , such that

$$P_j = Q_j P_\lambda^+ T_1 (P_\lambda^+ L_1)^k, \tag{89}$$

where  $Q_j$  is bounded and uniformly on  $\Gamma_\varepsilon$ . Since

$$\begin{aligned} & \| Q_j P_\lambda^+ T_1 (P_\lambda^+ L_1)^k D_\lambda^+ K \| \\ & \leq \| Q_j P_\lambda^+ \| \cdot \| T_1 (P_\lambda^+ L_1)^k D_\lambda^+ K \|, \end{aligned} \tag{90}$$

so, it is sufficient to prove

$$|\Im \lambda| \| T_1 (P_\lambda^+ L_1)^k D_\lambda^+ K \|, \tag{91}$$

is bounded and uniformly on  $\Gamma_\varepsilon$ . Since  $J_1$  is compact, it is sufficient to establish the result for an operator of rank one, that is  $T_1 : \varphi(a, v, \mu) \rightarrow T_1 \varphi(-a, v, \mu)$ ,

$$\begin{aligned} T_1 \varphi(-a, v, \mu) &= \theta(x) \int_E dv f(v, \mu) \\ & \times \int_0^1 g(v', \mu') \varphi(a, v', \mu') |\mu'| d\mu', \end{aligned} \tag{92}$$

where  $f(\cdot, \cdot), g(\cdot, \cdot)$  are measurable simple functions. For  $\varphi \in X$ ,

$$\begin{aligned} & T_1 (P_\lambda^+ L_1)^k D_\lambda^+ K \varphi(x, v, \mu) \\ &= \theta(x) f(v, \mu) \int_E dv' \int_0^1 d\mu' g(v', \mu') \\ & \times f(v', \mu') \int_{-a}^a dx \theta(x') \exp \left[ \frac{-1}{\mu'} \int_{x'}^x \right. \\ & \quad \left. \times (\lambda - \sigma(\xi, v)) d\xi' ((2k+1)a - x) \right] \int_E dv'' \\ & \times \int_{-1}^1 g(v'', \mu'') \varphi(x - s\mu', v'', \mu'') d\mu''. \end{aligned}$$

(93)

Now, we define the operators by

$$B_1 : \varphi \in X \rightarrow \theta(x) \int_E dv \times \int_{-1}^1 g(v, \mu) \varphi(x, v, \mu) d\mu, \quad (94)$$

$$B_2 : \gamma \in R \rightarrow \gamma \eta(v, \mu) \in L_{p,1}^i, \quad (95)$$

$$B_k : \varphi \in L_p(-a, a) \rightarrow \int_E dv \int_0^1 d\mu \times f(v, \mu) g(v, \mu) \int_{-a}^a \exp \left[ \frac{(2k+1)a-x}{\mu} \right] \times \theta(x) \int_{x'}^x (\lambda - \sigma(\xi, v)) d\xi \varphi(x) dx, \quad (96)$$

so, we can get

$$T_1(P_\lambda^+ L_1)^k D_\lambda^+ K = B_2 B_k B_1.$$

Clearly,

$$\| T_1(P_\lambda^+ L_1)^k D_\lambda^+ K \| \leq \| B_2 \| \cdot \| B_k \| \cdot \| B_1 \|,$$

because of  $B_1, B_2,$  and  $B_k$  are all bounded, moreover  $B_1$  and  $B_2$  are independent of  $\lambda$ , we only need to prove that  $\| \mathfrak{S}\lambda \| \| B_k \|$  is bounded uniformly on  $\Gamma_\varepsilon$ . Now, we set  $\varphi \in L_p((-a, a); dx)$  and  $\bar{\varphi}$  denote by its trivial extension to  $R$ , so  $B_k \varphi$  may be written in the from

$$B_k \varphi = \int_R F_\lambda((2k+1)a-x) \bar{\varphi} dx = (F_\lambda * \bar{\varphi})((2k+1)a).$$

the Young inequality gives

$$\| B_k \varphi \| \leq \| F_\lambda \|_{L^q(R)} \cdot \| \bar{\varphi} \|_{L^p(-a,a)},$$

then

$$\| B_k \| \leq \theta(x) \int_0^{+\infty} \left| \int_E dv \times \int_0^1 d\mu f(v, \mu) g(v, \mu) \exp \left[ -\frac{-1}{\mu} \right] \times \int_{x'}^x (\lambda - \sigma(\xi, v)) d\xi \right| d\mu dt.$$

Since

$$\| \mathfrak{S}\lambda \|^q \int_\varepsilon^{+\infty} \left| \int_E dv \int_0^1 d\mu f(v, \mu) \times g(v, \mu) \theta(x) \exp \left[ -\frac{1}{\mu} \int_{x'}^x \right] \times (\lambda - \sigma(\xi, v)) d\xi \right|^q dt,$$

is bounded uniformly on  $\Gamma_\varepsilon$ , we can get the (87) is bounded uniformly on  $\Gamma_\varepsilon$ .

Now, we prove that (88) is bounded uniformly on  $\Gamma_\varepsilon$ . Since  $H_{12} = \alpha T_1 + \beta L_1, T_1$  is compact operator, it suffices to prove that  $\| \mathfrak{S}\lambda \| \| K Q_\lambda^+ L_1 P_{2^n} D_\lambda^+ K \|$  is bounded uniformly on  $\Gamma_\varepsilon$ .

In fact, for all  $\varphi \in X$ , then

$$K Q_\lambda^+ L_1 (\beta P_\lambda^+ L_1) D_\lambda^+ K = \theta(x) \beta^n f(v, \mu) \int_{-a}^a dx \int_E dv' \int_0^1 g(v', \mu') \times f(v', \mu') \exp \left[ \frac{-(2n+1)a+x}{\mu'} \right] \times \int_{x'}^x (\lambda - \sigma(\xi', v)) d\xi' \times \theta(x') \int_E dv'' \int_{-1}^1 g(v'', \mu'') \times \varphi(x - s\mu', v'', \mu'') d\mu''.$$

Setting

$$K Q_\lambda^+ L_1 (\beta P_\lambda^+ L_1) D_\lambda^+ K = E_2 E_n E_1,$$

where

$$E_1 : \varphi \in X \rightarrow \theta(x) \int_E dv \int_{-1}^1 g(v, \mu) \times \varphi(x, v, \mu) d\mu \in L_p((-a, a); dx),$$

$$E_2 : \varphi \in L_p((-a, a); dx) \rightarrow \theta(x) \beta^n f(v, \mu) \varphi(x) \in X_p,$$

$$E_n : \varphi \in L_p((-a, a); dx) \rightarrow \int_{-a}^a dx \int_E dv \int_0^1 f(v, \mu) g(v, \mu) \times \exp \left[ \frac{1}{\mu} \int_{x'}^x (-\lambda + \sigma(\xi, v)) d\xi \right] \times ((2n+1)a-x) d\mu \in L_p((-a, a); dx).$$

Since operator  $E_1, E_2$  and  $E_n$  are bounded, moreover  $E_1$  and  $E_2$  are independent of  $\lambda$ , so it is easy to prove that,  $\| \mathfrak{S}\lambda \| \| E_n \|$  is bounded uniformly on  $\Gamma_\varepsilon$ . This ends the step three.

Finally, since (63) and equation (62) have the same form, so we can get the (63) is bounded uniformly on  $\Gamma_\varepsilon$ . This ends the proof.  $\square$

**Theorem 4.** *If assumption  $O_1$  and  $O_2$  are satisfied, then for all  $\varepsilon > 0$ , and big enough  $|\mathfrak{S}\lambda|$ , then the spectrum of transport operator  $A_H$  consists of, only, finite isolated eigenvalues which have a finite algebraic multiplicities in trip  $\Gamma_\varepsilon$ .*



**Proof.** On one hand, because of hypothesis  $O_2$  and Lemma 1 we can get operators  $K(\lambda I - B_H)^{-1}$  and  $(\lambda I - B_H)^{-1}K$  are compact operator on  $X$ . So, for all  $\lambda \in \Gamma_\varepsilon$ , the operator  $[(\lambda I - B_H)^{-1}K]^2$  is compact operator on  $X$ .

On the other hand, since

$$\begin{aligned} & [(\lambda I - B_H)^{-1}K]^2 \\ &= (\lambda I - B_H)^{-1}[K(\lambda I - B_H)^{-1}K], \end{aligned}$$

so, according to the Lemma 1 and Theorem 3, we can get

$$\lim_{|\Im \lambda| \rightarrow +\infty} \| [(\lambda I - B_H)^{-1}K]^2 \| = 0.$$

Finally, according to the Lemma 2, the desired result follows.  $\square$

**Acknowledgements** The authors would like to express their gratitude to the referee for his valuable suggestions. The research was supported by the National Natural Science Foundation of China (No.11461055) and the Natural Science Foundation of Jiangxi Province of China (No.20132BAB201002).

#### References:

- [1] J. Lehner and G. M. Wing, On the spectrum of an unsymmetric operator arising in the transport theory of neutron, *Communications on Pure and Applied Mathematics*, 8, 1955, pp.217-234 .
- [2] K. Latrach and A. Dehici, Spectral properties and time asymptotic behaviour of linear transport equations in slab geometry, *Mathematical Methods in the Applied Sciences*, 24, 2001, pp.689-711 .
- [3] K. Latrach, Description of the real point spectrum for a class of neutron transport operators, *Transport Theory and Statistical Physics*, 22(5),1993, pp.593-629.
- [4] K. Latrach and J. Martin Paoli, Relatively compact-like perturbations, *Journal of the Australian Mathematical Society*, 77, 2004, pp.73-89.
- [5] M. Mokhtar-Kharroubi, Time asymptotic behaviour and compactness in transport theory, *European Journal of Mechanics B:Fluids*,11, 1992, pp.39-68 .
- [6] S. H. Wang and Y. G. Zheng, The spectrum of transport operator with general boundary condition in slab geometry, *Mathematics in Practice and Theory*, 36(10),2006, pp.147-153.
- [7] S. H. Wang, M. Z. Yang and G. Q. Xu, The spectrum of the transport operator with generalized boundary conditions, *Transport Theory and Statistical Physics*, 25(7),1996, pp.811-823.
- [8] K. Latrach and A. Jeribi, On the essential spectrum of transport operators on  $L_1$  spaces, *Journal of Mathematical Physics*, 37(12),1996, pp.6486-6494.
- [9] M. Sbihi, A resolvent approach to the stability of essential and critical spectra of perturbed  $C_0$ -semigroups on Hilbert spaces with applications to transport theory, *Journal of Evolution Equations*, 7, 2007, pp.35-58.
- [10] M. Boulanouar, The asymptotic behavior for the transport operator in slab geometry. *Transport Theory and Statistical Physics*, 37, 2008, pp.38-64.
- [11] M. Boulanouar, Generation theorem for the streaming operator in slab geometry, *Journal of Dynamical and Control Systems*, 9, 2003, pp.33-51.
- [12] J. Voigt and B. Virginia, A perturbation theorem for the essential spectral radius of strongly continuous semigroups, *Mh. Math.*, 90,1980 , pp.153-161.
- [13] S. H. Wang and J. S. Ma, The spectrum of a singular transport operator with periodic boundary conditions in slab geometry, *Mathematics in Practice and Theory*, 37(6), 2007, pp.160-165.
- [14] S. H. Wang and H. X. Wu, The spectral analysis of transport operator with reflecting boundary condition in slab geometry, *Mathematics in Practice and Theory*, 38(22), 2008, pp.197-203.
- [15] K. Latrach and H. Megdiche, Spectral properties and regularity of solutions to transport equations in slab geometry, *Mathematical Methods in the Applied Sciences*, 29, 2006 , pp.2089-2121.
- [16] K. Latrach, H. Megdiche and M. A. Taoudi. Compactness properties for perturbed semigroups in Banach spaces and application to a transport model, *Journal of Mathematical Analysis and Applications*,359, 2009, pp.88-94.
- [17] H. X. Wu, S. H. Wang and D. B. Yuan, The spectrum distribution of transport operator with abstract boundary conditions, *Collectanea Mathematica Pure and Applied Mathematics*, vol.29, no.5, pp.489-497, 2013.
- [18] K. Latrach, On the spectrum of the transport operator with abstract boundary conditions in slab geometry, *Journal of Mathematical Analysis and Applications* , 252, 2000, pp.1-17.
- [19] H. X. Wu, S. H. Wang and D. B. Yuan, Spectral distribution of transport operator arising in cell population, *Journal of Function Spaces*, 2014, Article ID 748792, 10 pages.

- [20] K. Latrach and B. Lods, Spectral analysis of transport equations with bounce-back boundary conditions, *Mathematical Methods in the Applied Sciences*, 32, 2009, pp.1325-1344.
- [21] S. H. Wang, G. F. Cheng and D. B. Yuan. Spectral analysis of transport operator in Lebowitz-Rubinow model, *WSEAS Transactions on Mathematics*, 13, 2014, pp.324-334.
- [22] M. Z. Yang and G. T. Zhu, Spectrum of neutrontransport operator with anisotropic scattering and fission, *Scientia Sinica*, 4, 1981, pp.476-482.
- [23] G. Q. Xu, M. Z. Yang and S. H. Wang, On the eigenfunction expansion of the transport Semigroup for a Bounded Convex Body, *Transport Theory Statistical Physics*, 26(2), 1997, pp. 271-278.
- [24] W. Z. Yuan and G. Q. Xu, Spectral analysis of a two unit deteriorating standby system with repair, *WSEAS Transactions on Mathematics*, 10(4), 2011, pp. 125-138.
- [25] A. Jeribi, S. Ould Ahmed Mahmoud, and R. Sfaxi, Time asymptotic behaviour for a one-velocity transport operator with Maxwell boundary condition, *Acta Applicandae Mathematicae* 3, 2007, pp. 163-179.
- [26] S. Char, A. Jeribi and N. Moalla Time asymptotic behavior of the solution of an abstract Cauchy problem given by a one-velocity transport operator with maxwell boundary condition, *Collectanea Mathematica*, 64, 2011, pp. 97-109.
- [27] M. Chabi and K. Latrach. Singular one-dimensional transport equations on  $L_1$  space, *Journal of Mathematical Analysis and Applications*, 383, 2003, pp. 319-336.
- [28] M. Mokhtar-Kharroubi, New generationtheorems in transporttheory, *Afrika. Matematika*. 22, 2011, pp. 153-176.
- [29] B. Abdelmoumen, A. Jerihi and M. Mnif, Time asymptotic description of the solution to an abstract Cauchy problem and application to transport equation, *Mathematische Zeitschrift*, 268, 2011, pp. 837-869.
- [30] B. Lods and M. Sbihi, Stability of the essential spectrum for 2D-transport models with Maxwell boundary conditions, *Math. Meth. Appl. Sci.*, 29, 2006, pp. 499-523.
- [31] L. X. Ma, G. Q. Xu, N. E. Mastorakis, Analysis of a deteriorating cold standby system with priority, *WSEAS Transactions on Mathematics*, vol.10, 2011,2: pp.84-94.
- [32] M. M. Xu, J. G. Jia, P. Zhao. The progress for Stability of essential and critical spectral of perturbed Co-semigroups and its applications to models of transport theory, *WSEAS Transactions on Mathematics*, vol.11, 2012,6:pp.557-566.
- [33] M. Boulanouar, Transport equation in cell population dynamics II, *Electronic journal of differential equation*, 145, 2010, pp.1-20.
- [34] J. L. Lebowitz and S. I. Rubinow, A theory for the age and generation time distribution of a microbial population, *J.Math.Biol*, 1974, pp.17-36.
- [35] B. Lods and M. Mokhtar-Kharroubi, On the theory of a growing cell population with zero minimum cycle length, *Journal of Mathematical Analysis and Application*, 266, 2002, pp.70-99.