Usage of Finite Grid Solution for Vibration of Rectangular Plates on Elastic Foundations

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Abstract: - In engineering practice, beside static case often dynamic effects must be taken into consideration for plate design problems. Plate vibration solutions have been available for regular geometries for a long time, but where irregular boundaries or partial contact are encountered difficulties arise because it will be necessary to describe the governing equation of motion in a general mathematical form. This is not easy. The intention of this study is to extend analytical solutions of the discrete one-dimensional elements resting on elastic foundation for solution of plate vibration problems. The solution can be stated as an extension of the so-called discrete parameter approach where the physical domain is broken down into discrete sub-domains, each endowed with a response suitable for the purpose of mimicking problem at hand. The numerical results of this solution technique show versatile and powerful analysis capacity to solve various vibration problems including bi-directionally stepped and plates on non-homogeneous elastic foundation.

Key-Words: Grillage of beams; vibration; plates with abrupt changes in thickness; complex support conditions; finite grid solution

1 Introduction

Plates on elastic foundations have received considerable attention due to their wide applicability many engineering disciplines. Since the in interaction between structural foundations and supporting soil has a great importance in many engineering applications, a considerable amount of research has been conducted. Many studies, such as [1-9] have been done to find a convenient representation of physical behavior of a real structural component supported on a foundation. There are several practical foundation models as well as their proper mathematical formulations. A broad range of the beam or plates as engineering problems has been solved numerically such as finite element and boundary element methods [10-15]. Owing to its convenience in solution of plate problems as a numerical method the finite strip method have attracted much attention from many authors as [16-18] suggested a procedure incorporating the finite strip method together with spring systems for treating plates on elastic supports. However series and closed form solutions

for plates have been published for a limited number of cases as [19-25]. The orthogonalization of the series and other calculations are performed using Fourier expansion of Bernoulli polynomials under some realistic approximations for the limiting values of the boundary conditions. The studies can be summarized as series expansion consisting of some specially chosen trigonometric functions used for free vibrations of rectangular plates resting on elastic foundations with various boundaries and subjected to uniform and constant compressive, unidirectional forces and closed form solutions of free vibration problem of thin rectangular plates on Winkler and Pasternak elastic foundation model developed with some limitations such as mixed or fully-clamped boundary conditions etc.

This study is oriented toward the development of finite grid elementas an application of the finite element method. The aim is to investigate an improved finite grid solution for vibration problems of plates on elastic foundation. In this method the discretized plate element is reassembled by the matrix displacement method so that consistent mass matrix of the total structure is generated schme to compute all displacements for each nodal point in a convenient sequence.

2 Theory of Problem Formulation

In engineering practice, dynamic effects need to be taken into consideration for a wide variety of plate problems. It will be necessary to describe the governing equation of motion of plates in a general mathematical form for such cases. This can be achieved by inserting the inertia forces due to the lateral translations, in an appropriate way, into the governing differential equation for static equilibrium. For dynamic problems of the plates on elastic foundations with arbitrary shapes and boundary conditions with most elements developed to date there exists no rigorous solution except in the form of infinite Fourier series for a Levy-type solution. The series solutions are valid for very limited cases such as when the second foundation parameter has been eliminated, and simple loading and boundary conditions exist. Grillages of beam elements that have no such limitations can represent the plates.



Fig.1 Representation of a rectangular plate by grids as parallel sets of one-dimensional beam elements replaces the continuous surface.

The usual approach in formulating problems of beams, plates, and shells supported by elastic media is based on the inclusion of the foundation reaction in the corresponding differential equation of the beam, plate, or shell. In case of elastic foundation under the combined action of transverse load and vibration the governing differential equation of the plates can be obtained as;

$$D(\frac{\partial^4 w}{\partial x^4} + 2\frac{\partial^4 w}{\partial x^2 \partial y^2} + \frac{\partial^4 w}{\partial y^4}) + k_I w$$

$$-k_{\theta}(\frac{\partial^2 w}{\partial x^2} + \frac{\partial^2 w}{\partial y^2}) - \overline{m}\frac{\partial^2 w}{\partial t^2} = q(x, y)$$
(1)

where w = w(x, y, t) is the transverse deflection of the plate, k1 is Winkler parameter with the unit of force per unit area/per unit length (force/length³), k₀ is reaction moment per unit area per unit rotation, q(x,y) is the external loads, D is flexural rigidity of plate and \overline{m} is the mass of the plate per unit area.

By representing the plate with assemblage of individual beam elements interconnected at their neighboring joints, the system cannot truly be equal to the continuous structure, however sufficient accuracy can be obtained similar to the static case. Therefore plates can be modeled as an assemblage of individual beam elements interconnected at their intersecting joints. There are many researches concerning analysis of beam element resting on elastic foundation as [26-30]. The properties of such beam elements on elastic foundations will be a very useful tool to solve such generalized problems.

By representing the plate shown in Fig.1 with individual beam elements the problem can be reduced to a one-dimensional one. On the other hand the similar elements can be formed in radial and tangential directions for circular plates as [31]. By representing the continuous plate with individual beam elements resting on continuous springs. The problem will be reduced to one-dimensional one. Then Eq. (1) can be rewritten in reduced form of the governing equation for one-dimensional beam elements as;

$$D\frac{d^4w}{dx^4} + k_1w - k_\theta \frac{d^2w}{dx^2} + \overline{m}\frac{d^2w}{dt^2} = 0$$
(2)

The main advantage of the reduction is that the exact geometric stiffness matrix can be determined for the beam elements and these matrices can be used as a basis of assembling the elements to apply to plate problems as [32]. Then dynamic problems of the plates resting on Winkler foundation with arbitrary loading and boundary conditions could be solved approximately. Assemblies of beam elements that have no limitations for loading and boundary conditions can represent plates adequately.

The degrees of freedom of the element are the local torsion, rotation and translation at each end. Since the angular displacements are obtained from the pure torsion member, the torsional DOF's are independent of the foundation. Then it can be assumed that the displacements within the span are defined by the same interpolation functions those already derived for obtaining the element stiffness matrices, loading conditions and discontinuous surfaces.

3 Consistent Mass Matrices

Consider the beam element shown in Fig. 2 having a mass distribution m(x). If it is subjected to a unit angular acceleration at point a, the acceleration would be developed along its length as follow;

$$\ddot{w}(x) = \psi_2(x)\ddot{w}_2 \tag{3a}$$

By d'Alembert's principle, the inertial force due to this acceleration is;

$$f_1(x) = m(x)\ddot{w}(x) = m(x)\psi_2(x)\ddot{w}_2$$
 (3b)





Fig.2 Representation of a beam element subjected to a unit real acceleration and virtual translation at the left side

By the principle of virtual displacements the mass influence coefficients associated with this acceleration as the nodal inertial forces can be evaluated. As an example, it is possible to evaluate the vertical force pa, equating work done by the external force due to virtual displacement, to the work done on the distributed inertial forces $f_I(x)$. That is,

$$p_a \delta w_3 = \int_0^L f_I(x) \delta w(x) dx$$
(4)

Substituting the vertical virtual displacement in terms of the shape functions into the equation then,

$$m_{23} = \int_{0}^{L} m(x) \psi_{2}(x) \psi_{3}(x) dx$$
 (5)

By this analogy, this equation can be extended to evaluate for the other degrees of freedoms such as;

$$m_{ij} = \int_{0}^{L} m(x) \psi_{i}(x) \psi_{j}(x) dx$$
 (6)

By using the proper shape functions, the corresponding shape functions derived for conventional beam or beam element resting one or two-parameter elastic foundations, this equation lets to evaluate all of the mass matrix terms. Computing the mass coefficients by the same shape functions with same procedures as done for determining the stiffness matrices is called consistent-mass matrices.

3.1 Consistent mass matrices for Two-Parameter foundation

For one-parameter foundation case it is possible to evaluate mass influence coefficients of a structural element with the procedures similar to that obtaining the element stiffness matrix by making the use of finite element concept. The consistent mass matrix of beam elements resting on two-parameter elastic foundations can also be evaluated by the same procedures as Winkler parameter case [33]. Substituting the proper shape functions of the beam elements resting on two-parameter derived by [32] $A < 2\sqrt{B}$ $A > 2\sqrt{B}$ and cases for both respectively, into Equation (6) leads to evaluate the consistent mass matrices. The terms of the mass matrix, $m_{ij} = f(p,t,\mu,L)$, obtained as functions of foundation parameters, length of the elements and mass per unit length. Since the terms for twoparameter cases are too complex and extremely long functions, they are not presented in this study. However, by letting both of the foundation parameters tend to zero, the correctness of the terms is checked. When foundation parameter k_1 and k_{θ} tend to zero (or $p \rightarrow 0$ and $t \rightarrow 0$), the terms in the equations reduce to the conventional beam consistent mass terms obtained by Hermitian functions as for Winkler case.

$$\underline{Lim}_{p\to 0}[\underline{M}] = \frac{\mu L}{420} \begin{vmatrix} 140 & 0 & 70 & 0 & 0 \\ 0 & 4L^2 & -22L & 0 & -3L^2 & -13L \\ 0 & -22L & 156 & 0 & 13L & 54 \\ 70 & 0 & 0 & 70 & 0 & 0 \\ 0 & -3L^2 & 13L & 0 & 4L^2 & 22L \\ 0 & -13L & 54 & 0 & 22L & 156 \end{vmatrix}$$

In two parameter case same new parameters introduced as

$$\alpha = \sqrt{\lambda^2 + \delta} = \lambda \sqrt{1 + t}$$

$$\beta = \sqrt{\lambda^2 - \delta} = \lambda \sqrt{1 - t}$$

and

$$\alpha = \sqrt{\lambda^2 + \delta} = \lambda \sqrt{1 + t}$$

$$\beta = \sqrt{\delta - \lambda^2} = \lambda \sqrt{t - 1}$$
 For $A > 2\sqrt{B}$

where t is dimensionless as;

$$T = \frac{\delta}{\lambda^2} = \frac{\frac{M_{\theta}}{4EI}}{\sqrt{\frac{k_1}{4EI}}}$$

k.

The influence of the foundation parameters k_1 and k_{θ} on the consistent mass terms for $A < 2\sqrt{B}$ with corresponding terms of Eq. (6) can be normalized as shown in Fig. 4. Note that, as the second parameter tends to zero (i.e. $t \rightarrow 0$) the same two-dimensional curves of one-parameter case given in Fig. 3 are obtained.



Fig.3 Influence of two-parameter foundation on the m_{22} , m_{25} , m_{26} , m_{33} , m_{36} and m_{56} normalized consistent mass terms

From the figure it is inferred that presence of second foundation parameter k_{θ} in the analysis is remarkably dominant. This might have been anticipated because strain energy density functional includes one more term in the case of two parameter foundation than that of the Winkler foundation.

3.2 Assembling the consistent mass matrix of the total structure

After obtaining the consistent mass matrices of each one dimensional elements the discretized plate element reassembled by the matrix displacement method to obtain free vibration frequencies of a total structure. That is, the stiffness and consistent mass matrices of the total structure is generated by using a proper numbering shame to collect all displacements for each nodal point in a convenient sequence of the system for rectangular grids can be generated as follow;

$$\underline{\boldsymbol{M}}_{sys} = \sum_{i=1}^{NE} \ \underline{\boldsymbol{a}}_{i}^{T} \underline{\boldsymbol{M}}_{i} \ \underline{\boldsymbol{a}}_{i}$$

where i is the individual element number, NE is the number of elements depending on boundary conditions, a_i is the individual rotation element matrix, M_i is the proper element consistent mass matrix for a beam conventional resting on oneparameter elastic foundation and M_{sys} is the system consistent mass matrix. Then the equations of motion for a system in a free vibration as an eigenvalue problem may be written as;

$$(\underline{k}_{svs} - \omega^2 \underline{M}_{svs}) \underline{w} = 0$$

where the quantities ω^2 are the eigenvalues indicting the square of free vibration frequencies that satisfy the above equation, while the corresponding displacement vector w express the fitting shapes of the vibrating system as the eigenvectors of mode shapes and k_{sys} is the stiffness matrix of the total structure defined by the same interpolation functions those already derived for obtaining the element stiffness matrices.

4 Case Study

The validity of the finite grid solution (FGM) were carried out with some comparisons of previous studies in literature for vibration problems of plates resting on elastic foundations [32]. The examples of the references were verified to demonstrate the accuracy of the method for vibration problems of plates on elastic foundations.

In this study the examples concern with the free vibration of a directionally stepped plate and a uniform stepped plate on non-homogeneous elastic foundations studied by [34] as examples of analysis of plates with abrupt changes in thickness and complex support conditions. The fundamental mode of the plate simply supported all edges with abrupt changes in thickness shown in Fig.4 studied to indicate convergence and accuracy of the method for variation of thickness ratio. The plate is considered to be isotropic with elasticity module E, Poisson's ratio v, flexural rigidity D of the thicker part and ρ is the mass density.



Fig.4 Bi-directionally stepped square plate with all edges simply supported (a) plan; (b) section

The fundamental angular frequency w_1 is expressed in terms of a dimensionless parameter β_1 can be presents as

$$w_1 = \frac{\beta_1}{4a^2} \sqrt{\frac{D}{\rho h_1}}$$
 where $D = \frac{Eh^3}{12(1-v^2)}$

The results of variation of fundamental angular frequency with the thickness ratio h_2/h_1 presented in Fig. 5 shows a good agreement between the methods.



Fig.5 Comparison of fundamental frequencies for bi-directionally stepped and simply supported (SSSS) square plate

For fundamental angular frequencies of a uniform square plate resting on a non-homogeneous foundation shown in Fig. 6 is considered to check the validity of the finite grid solution. The plate is supported on elastic foundation of modulus K2 within the central square region of size 1.2ax1.2a, and elsewhere the foundation modulus is K1. The fundamental angular frequencies are compared for all of the edges simple supported (SSSS) boundary conditions.



Fig.6 A uniform square plate on non-homogeneous elastic foundation (a) plan, (b) section

The fundamental angular frequencies are compared for all of the edges simple supported (SSSS) boundary conditions as shown in Fig. 7.



Fig.7 Fundamental frequency coefficients of square plate on non-homogeneous elastic foundation (SSSS)

The results for the plate with all edges simple supported portrayed graphically in Figs. 5 and 7. The solution method shows good agreement observed with the reference values for all cases as illustrated for solution of vibration problems.

However in order to present a comparison study for vibration of plates resting on 2-parapeter foundation case free vibration analyses of square Levy-plates studied by Lam et al [35] taken into consideration.

For the square plates with different combinations f or values of the non-dimensional foundation parameters k_1 and k_2 defined in Eq. (14), the fundamental frequencies are compared with the reference results for various foundation parameters as shown in Fig. 8. As for the other cases the vibration of plates on two-parameter foundation results show that the method produces very good results for the computational effort that goes into it.



Fig.8 Fundamental frequency parameter of the four edges simple supported square plate on two-parameter foundation (SSSS and v = 0.3)

5 Conclusion

The solution of free vibration problems for rectangular plates resting on elastic foundations is considered to be too complex. In many cases there is apparently no analytical solution other than simple cases. A grid work analogy called the Finite Grid Solution involving discretized plate properties mapped onto equivalent beams with adjusted parameters and matrix displacement analysis are used to develop a more general simplified numerical approach for such complicated problems. It is shown that after obtaining solutions of the governing differential equations of beam elements, the derived exact shape functions (interpolation functions) have extended to determine consistent mass matrices by finite element method.

It is noted that the consistent mass terms related to one dimensional beam elements on elastic foundations are very sensitive to variation of foundation parameters. It can be concluded that the finite grid solution as a combination of finite element method, lattice analogy and matrix displacement analysis of grid works is a useful tool to improve the solution for various vibration of plate problems. The numerical results of various vibration problems of plates on elastic foundations analysis of plates with abrupt changes in thickness and complex support conditions show that the finite grid solution is versatile and has high accuracy. On the other hand it is possible to extend the method to solve more complex cases of plate, grid and shell problems for various boundary conditions as an approximate solution. This is possible for free as well as forced vibration cases.

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